

Compressible MHD Turbulence: A Source of Heating in the Fast Solar Wind

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June 28, 2016

Anomalous temperature profile of the solar wind

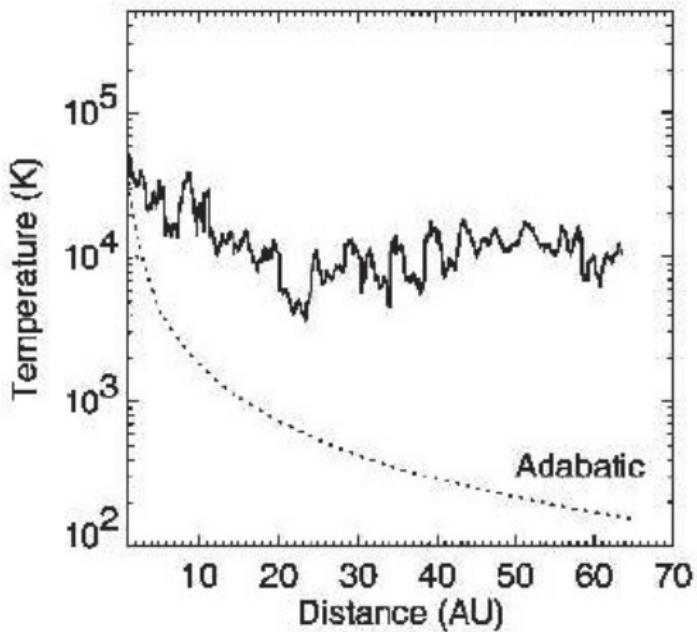


Figure : $T(r)$ of the solar wind (Voyager data, Richardson & Smith, GRL, 2003).

Required heating energy for the fast solar wind

- Power law for temperature profile of the FSW with distances
→ $T(r) \sim r^{-\xi}$
- Model relating ξ and heating energy flux rate ε_h (Vasquez *et al.*, JGR, 2007)

$$\varepsilon_h = \frac{3}{2} \left(\frac{4}{3} - \xi \right) \frac{V_{sw} k_B T(r)}{m_p r},$$

where V_{sw} be the solar wind speed, k_B the Boltzmann constant, T the temperature, m_p the protonic mass and r is the radial distance of the spacecraft from sun.

- At 1 AU, for $\xi = 0.72$ (from Voyager data) ⇒ required heating energy flux rate per unit volume as $5 \times 10^{-17} \text{ J.m}^{-3}.\text{s}^{-1}$

Turbulence as a heating source

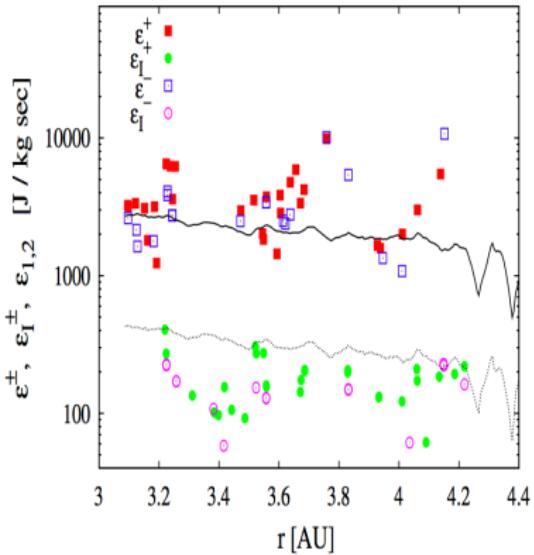


Figure : Turbulent contribution to SW heating (Carbone et al., PRL, 2009)

Incompressible MHD turbulence
(Politano & Pouquet, PRE, 1998)

$$\langle (\delta \mathbf{Z}^{\pm})^2 \delta Z_r^{\mp} \rangle = \frac{4}{3} \varepsilon_I^{\pm} r,$$

$$\text{where } \mathbf{Z}^{\pm} = \left(\mathbf{v} \pm \frac{\mathbf{b}}{\sqrt{\mu_0 \rho}} \right)$$

Heuristic compressible model
(Carbone et al., PRL, 2009)

$$\langle (\delta \mathbf{W}^{\pm})^2 \delta W_r^{\mp} \rangle = \frac{4}{3} \langle \rho \rangle \varepsilon^{\pm} r,$$

$$\text{where } \mathbf{W}^{\pm} = \rho^{1/3} \left(\mathbf{v} \pm \frac{\mathbf{b}}{\sqrt{\mu_0 \rho}} \right)$$

- The previous law is heuristic and in this case the model does not carry significant physical meaning.
- The frequency range chosen does not exactly correspond to usual MHD range ($10^{-4} - 10^{-1}$ Hz.).
- No intervals are found for which both of ε^\pm are conserved simultaneously and in fact
- In compressible MHD turbulence, the Pseudo energies $\frac{1}{2}\mathbf{W}^\pm \cdot \mathbf{W}^\pm$ are no longer inviscid invariants.

Can we do better ?

Exact relation for space plasma turbulence

Isothermal MHD turbulence (Banerjee & Galtier, PRE, 2013)

$$\begin{aligned} -2\varepsilon = & \frac{1}{2} \nabla_E \cdot \overbrace{\left\langle \left[\frac{1}{2} \delta(\rho \mathbf{z}^-) \cdot \delta \mathbf{z}^- + \delta \rho \delta e \right] \delta \mathbf{z}^+ + \left[\frac{1}{2} \delta(\rho \mathbf{z}^+) \cdot \delta \mathbf{z}^+ + \delta \rho \delta e \right] \delta \mathbf{z}^- + \bar{\delta}(e + \frac{v_A^2}{2}) \delta(\rho \mathbf{z}^- + \rho \mathbf{z}^+) \right\rangle}^{\text{Usual flux term}} \\ & - \underbrace{\frac{1}{4} \left\langle \frac{1}{\beta'} \nabla' \cdot (\rho \mathbf{z}^+ e') + \frac{1}{\beta} \nabla \cdot (\rho' \mathbf{z}'^+ e) + \frac{1}{\beta'} \nabla' \cdot (\rho \mathbf{z}^- e') + \frac{1}{\beta} \nabla \cdot (\rho' \mathbf{z}'^- e) \right\rangle}_{\text{New type of flux term}} \\ & + \left\langle (\nabla \cdot \mathbf{v}) \left[R'_E - E' - \frac{\bar{\delta}\rho}{2} (\mathbf{v}_A' \cdot \mathbf{v}_A) - \frac{P'}{2} + \frac{P'_M}{2} \right] \right\rangle + \left\langle (\nabla' \cdot \mathbf{v}') \left[R_E - E - \frac{\bar{\delta}\rho}{2} (\mathbf{v}_A \cdot \mathbf{v}_A') - \frac{P}{2} + \frac{P_M}{2} \right] \right\rangle \\ & + \left\langle (\nabla \cdot \mathbf{v}_A) [R_H - R'_H + H' - \bar{\delta}\rho(\mathbf{v}' \cdot \mathbf{v}_A)] \right\rangle + \left\langle (\nabla' \cdot \mathbf{v}_A') [R'_H - R_H + H - \bar{\delta}\rho(\mathbf{v} \cdot \mathbf{v}_A')] \right\rangle \end{aligned}$$

where

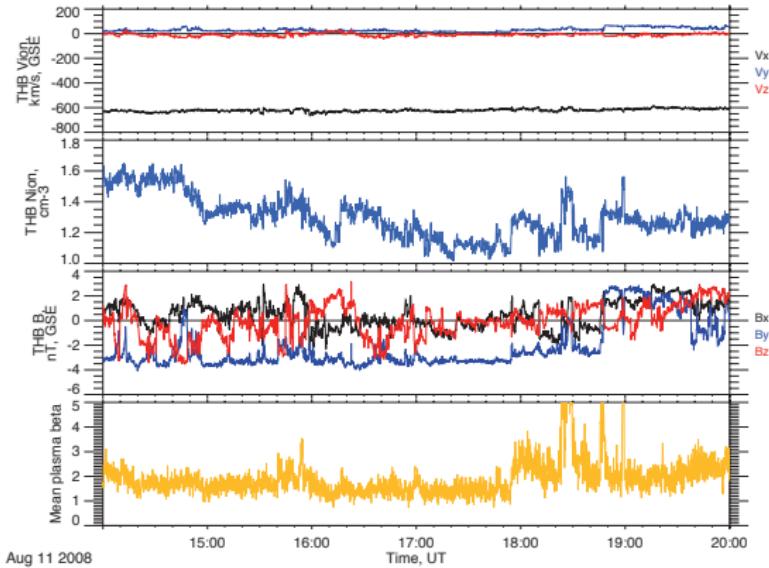
$$E = \rho(\mathbf{v} \cdot \mathbf{v} + \mathbf{v}_A \cdot \mathbf{v}_A)/2 + \rho e, \quad E' = \rho'(\mathbf{v}' \cdot \mathbf{v}' + \mathbf{v}'_A \cdot \mathbf{v}'_A)/2 + \rho' e';$$

$$R_E = \rho(\mathbf{v} \cdot \mathbf{v}' + \mathbf{v}_A \cdot \mathbf{v}'_A)/2 + \rho e', \quad R'_E = (\rho' \mathbf{v}' \cdot \mathbf{v} + \mathbf{v}'_A \cdot \mathbf{v}_A)/2 + \rho' e;$$

$$R_H = \rho(\mathbf{v} \cdot \mathbf{v}'_A + \mathbf{v}' \cdot \mathbf{v}_A)/2, \quad R'_H = \rho'(\mathbf{v}' \cdot \mathbf{v}_A + \mathbf{v} \cdot \mathbf{v}'_A)/2$$

$$H = \rho \mathbf{v} \cdot \mathbf{v}_A, \quad H' = \rho' \mathbf{v}' \cdot \mathbf{v}'_A; \quad \beta = 2C_S^2/v_A^2; \quad \beta' = 2{C_S}'^2/{v_A}'^2$$

Exploitation of compressible MHD turbulence model in solar wind



Solar wind data source: Themis B (NASA)

Plasma data : Electrostatic analyzer (ESA) – resolution 3 s.

magnetic data: Fluxgate Magnetometer (FGM) – resolution 3 s.

Compressible scaling in the solar wind

Schematic expression of newly derived exact relation

$$-2\varepsilon = \frac{1}{2} \nabla_{\mathbf{r}} \cdot (\langle \mathcal{F}_1 + \mathcal{F}_2 + \mathcal{F}_3 \rangle + \langle \Phi \rangle) + \text{source terms}$$

where

$$\langle \mathcal{F}_1 \rangle = \left\langle \left[\frac{1}{2} \delta(\rho \mathbf{z}^-) \cdot \delta \mathbf{z}^- \right] \delta \mathbf{z}^+ + \left[\frac{1}{2} \delta(\rho \mathbf{z}^+) \cdot \delta \mathbf{z}^+ \right] \delta \mathbf{z}^- \right\rangle,$$

$$\langle \mathcal{F}_2 \rangle = \langle 2\delta\rho\delta e\delta \mathbf{v} \rangle,$$

$$\langle \mathcal{F}_3 \rangle = \left\langle 2\bar{\delta}(e + \frac{v_A^2}{2})\delta(\rho \mathbf{v}) \right\rangle,$$

$$\langle \Phi \rangle = -\frac{1}{2} \left\langle \frac{1}{\beta'} \nabla' \cdot (\rho \mathbf{v} e') + \frac{1}{\beta} \nabla \cdot (\rho' \mathbf{v}' e) \right\rangle.$$

Intervals of uniform plasma beta

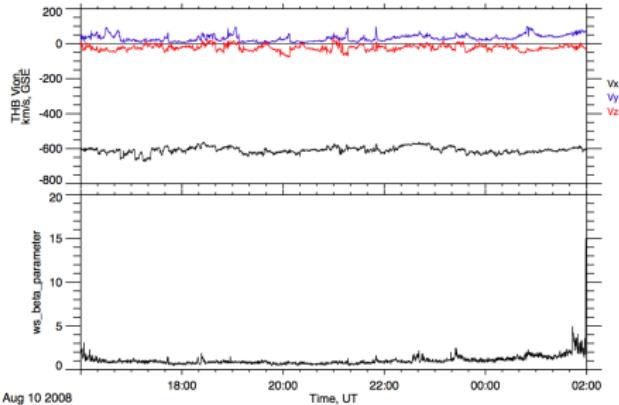


Figure : 10 hours interval with uniform β ($\simeq 1$)

For the intervals with $\beta \simeq 1$

$$\langle \Phi \rangle \approx -\frac{1}{2} \nabla_{\mathbf{r}} \cdot \langle \rho \mathbf{v} \mathbf{e}' - \rho' \mathbf{v}' \mathbf{e} \rangle = \nabla_{\mathbf{r}} \cdot \langle \bar{\delta} e \delta(\rho \mathbf{v}) \rangle$$

Heating energy from compressible MHD turbulence

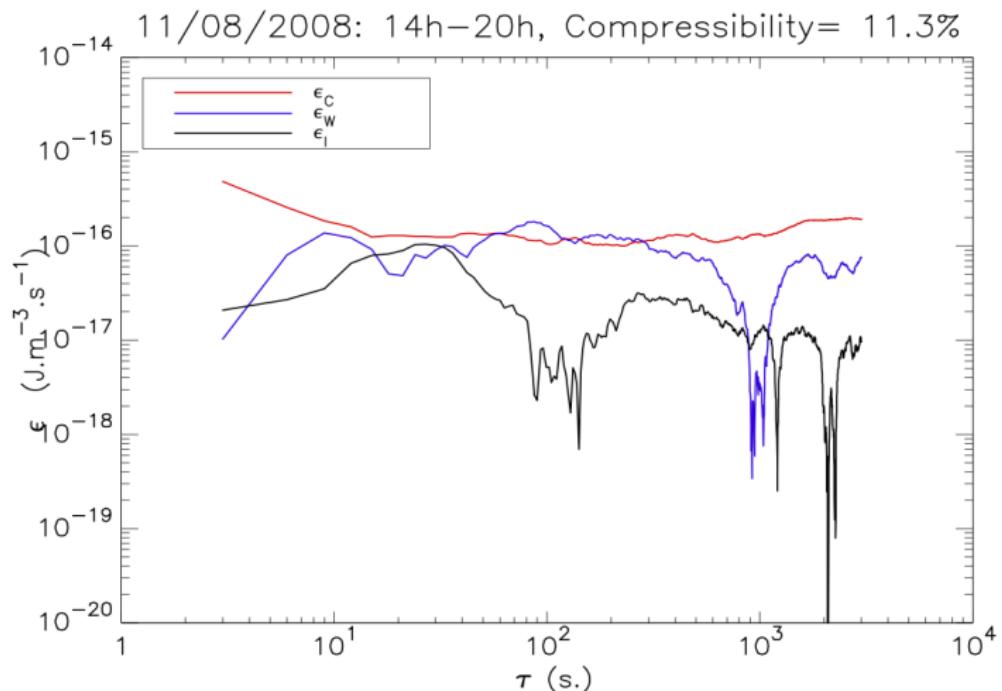


Figure : Estimates of energy cascade rate of different models

Comparison of different terms

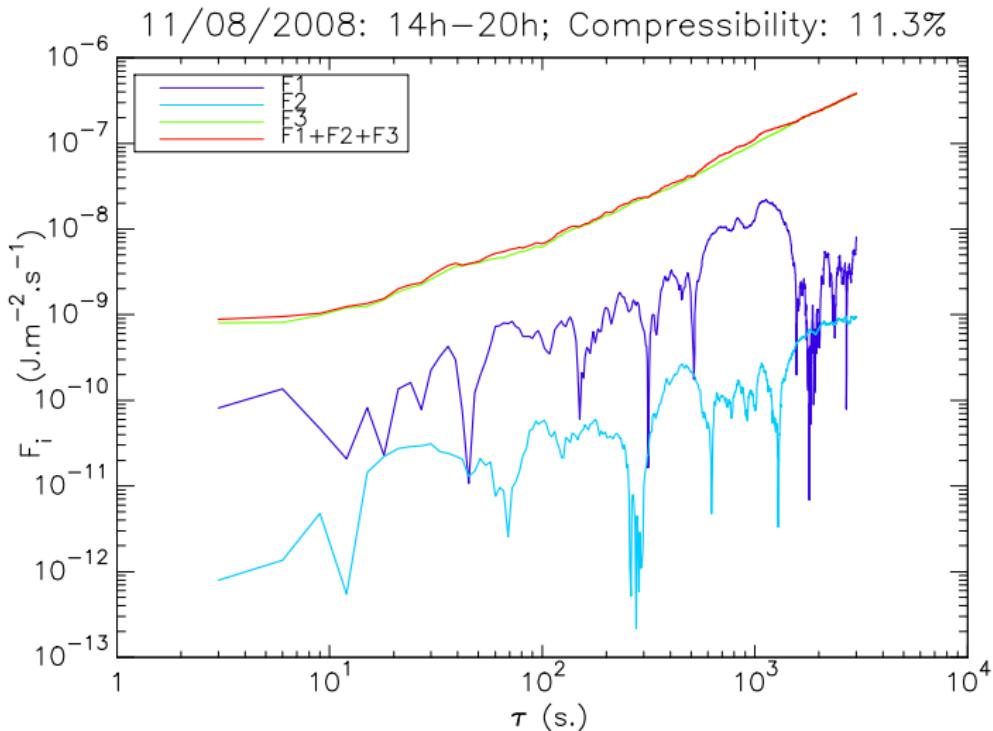


Figure : Comparison of the different terms F_1 , F_2 and F_3 .

Thank You